Witnessing the emergence of a carbon star

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ABSTRACT
During the late stages of their evolution, Sun-like stars bring the products of nuclear burning to the surface. Most of the carbon in the Universe is believed to originate from stars with masses up to a few solar masses. Although there is a chemical dichotomy between oxygen-rich and carbon-rich evolved stars, the dredge-up itself has never been directly observed. In the last three decades, however, a few stars have been shown to display both carbon- and oxygen-rich material in their circumstellar envelopes. Two models have been proposed to explain this dual chemistry: one postulates that a recent dredge-up of carbon produced by nucleosynthesis inside the star during the Asymptotic Giant Branch changed the surface chemistry of the star. The other model postulates that oxygen-rich material exists in stable keplerian rotation around the central star. The two models make contradictory, testable, predictions on the location of the oxygen-rich material, either located further from the star than the carbon-rich gas, or very close to the star in a stable disc. Using the Faint Object InfraRed CAmera (FORCAST) instrument on board the Stratospheric Observatory for Infrared Astronomy (SOFIA) Telescope, we obtained images of the carbon-rich planetary nebula BD+30 3639 which trace both carbon-rich polycyclic aromatic hydrocarbons and oxygen-rich silicate dust. With the superior spectral coverage of SOFIA, and using a 3D photoionization and dust radiative transfer model we prove that the O-rich material is distributed in a shell in the outer parts of the nebula, while the C-rich material is located in the inner parts of the nebula. These observations combined with the model, suggest a recent change in stellar surface composition for the double chemistry in this object. This is evidence for dredge-up occurring ∼10^3 yr ago.

Key words: circumstellar matter – ISM: abundances – planetary nebulae: individual: (BD+30 36 39) – infrared: stars.

1 INTRODUCTION
Planetary Nebulae (PNe) are the final evolutionary phase of low- and intermediate-mass stars. The nebulae form out of the mass lost by the star on the asymptotic giant branch (AGB), which may exceed 50 per cent of the stellar mass. The ejecta mainly consist of gas, initially molecular and atomic but becoming ionized by the remnant white dwarf star. Some solid particles (‘dust’) condense in the ejecta. The ejecta quickly disperses and merges with the surrounding interstellar medium. This recycled gas fuels most of the star formation in late-type galaxies (Leitner & Kravtsov 2011).

The chemical evolution of the Universe is driven by products of nucleosynthesis included in stellar ejecta. Low-mass stars are the primary source of new dust in the Milky Way (Matsuura et al.
2 OBSERVATIONS

Using the Faint Object InfraRed CAmera (FORCAST) for the Stratospheric Observatory for Infrared Astronomy (SOFIA) Telescope (Adams et al. 2010) we obtained images of BD +30° 3639. The observations cover the wavelength range of both the PAHs and crystalline silicate features. SOFIA uniquely allows for the imaging of the long wavelength emission from crystalline silicate features at high-angular resolution.

The FORCAST instrument has a short wavelength camera (SWC) that operates from 5–25 µm and a long wavelength camera (LWC) that operates from 25–40 µm, with several filters available in both cameras. FORCAST samples at 0.75 arcsec pixel−1, giving a 3.2 × 3.2 arcmin instantaneous field-of-view. Five different filters were used: FOR_F064 (centred at 6.4 µm with a Δλ = 0.14 µm), FOR_F077 (centred at 7.7 µm with a Δλ = 0.47 µm), FOR_F111 (centred at 11.1 µm with a Δλ = 0.95 µm), FOR_F113 (centred at 11.3 µm with a Δλ = 0.24 µm), and FOR_F336 (centred at 33.6 µm with a Δλ = 1.9 µm). We used the symmetric nod-match-chop (NMC) imaging mode, the chop is symmetric about the optical axis of the telescope with one of the two chop positions centred on the target. The nod throw is orientated 180° from the chop, such that when the telescope nods, the source is located in the opposite chop position.

3 RESULTS

In Fig. 1, we present the images of BD +30° 3639 obtained using the filters centred at 6.4, 7.7, 11.1, 11.3, and 33.6 µm. Each image has an overplotted blue circle of 5 arcsec radius in all of them. The white line in the top-left panel represents the PA of the cut made to all the images.

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Figure 2. Radial cuts along the major axis of BD +30° 3639, along a PA of 135° for all the SOFIA images: Top image shows the PAHs filters only (6.4, 7.7, 11.1, and 11.3 μm). The middle image shows the PAHs filters and the silicates filter (33.6 μm), while the bottom image shows all the filters with the flux normalized.

Figure 3. ISO spectrum of BD +30° 3639. The dashed lines are used to show the position and width of the SOFIA filters used on this Letter. Dark-blue represents the bandpass for the FOR_F064 filter, the FOR_F077 filter bandpass is in purple, brown for the FOR_F111 filter, green for the FOR_F113 filter, and light-blue for the FOR_F336 filter.

Extended than the C-rich material, with a FWHM of 12 arcsec. The bottom plot shows the same radial cuts done in all the filters with the normalized flux.

Fig. 3 shows the Infrared Space Observatory (ISO) Short Wave-length Spectrometer (SWS) spectrum of BD +30° 3639 (produced from Sloan et al. 2003), with the bandpass of the SOFIA filters overlaid by a dashed line of different colours depending on the filter. Dark-blue represents the bandpass for the FOR_F064 filter, the FOR_F077 filter bandpass is in purple, brown for the FOR_F111 filter, green for the FOR_F113 filter, and light-blue for the FOR_F336 filter. Using the ISO spectrum we can clearly see the PAHs and the silicates features covered by each filter. We can see the main PAHs (6.25, 7.7, and 11.3 μm) as well as the crystalline silicates (33.6 μm) features.

Previous observations of BD +30° 3639 show a shell-like structure with a radius of 5 arcsec x 4 arcsec in the visible (Li et al. 2002) and 6 arcsec x 5 arcsec in radio (Bryce et al. 1997), while a much larger extension of the faintest part of the nebula is detected in the near-infrared (Phillips & Ramos-Larios 2007). Previous mid-infrared imaging observations of BD +30° 3639 with a spatial resolution of 1.5 arcsec found that the 8.6 and 11.3 μm (PAHs) bands are slightly more extended than the continuum emission (Bentley et al. 1984). Higher spatial resolution imaging and spectroscopic observations of BD +30° 3639, taken with the Subaru telescope, confirmed that the PAH emission bands are more extended than the dust continuum emission, but show a similar extent to the [Ne II] 12.8 μm emission (Matsumoto et al. 2008). These observations also show evidence for a 10-μm absorption feature, attributed to silicates.

Until SOFIA, it was not possible to map the silicate emission features in this object, because the features emit at longer infrared wavelengths where no instrument was able to offer the necessary combination of wavelength coverage, sensitivity, and spatial resolution.Using the expansion velocity of the silicate shell of 10 km s^{-1}, and a distance of 1.2 kpc (Li et al. 2002), the expansion rate of the shell is 1.76 ± 0.25 mas yr^{-1} in radius. Using the images we compared the size of the C-rich to the O-rich shell. For the C-rich shell we averaged all the images of the PAHs filters (6.4 to 11.3 μm), for the O-rich shell we used the image of the silicates filter (33.6 μm). We measure the size (radius) of the C-rich and the O-rich shells to be 5 ± 0.75 arcsec and 7.5 ± 0.75 arcsec, respectively. These correspond to dynamical ages of 2800 ± 580 yr for the C-rich shell and 4300 ± 740 yr for its O-rich counterpart. This means that the transition of the star from O- to C-rich took place within a window of 1500 ± 940 yr. These values are entirely consistent with the post-AGB age of the central star determined by comparison of its luminosity and temperature to evolutionary tracks (Vassiliadis & Wood 1993; Blöcker 2001).
4 ANALYSIS AND DISCUSSION

In Figs 1 and 2, we can see the different extensions of the two major components presented here, PAHs and silicates. It is important to note that no continuum filter observations were made, as such, the dust continuum contribution is not subtracted from any of the presented observations. Therefore, one must consider the level of continuum contamination in each filter before drawing conclusions as to the relative distribution of PAHs and silicates.

In the case of PAHs, from Fig. 3, we can see that the \textit{FOR}_F064 filter has a negligible contamination by dust continuum and thus traces the PAHs emission. From Fig. 2 we can see that the emission in this filter has the same extension as the other PAHs filters, as such, we can assume that the measurements made in the PAHs filters are tracing the PAH emission, rather than just extended dust continuum.

Unfortunately, the observation made in the filter encompassing the silicates emission, the \textit{FOR}_F336 filter, includes a significant contribution of dust continuum, as shown in Fig. 3. Based on the images alone, we are unable to determine the spatial distribution of the silicates with respect to the dust continuum in the nebula. In order to resolve this degeneracy, a radiative transfer model of the nebula was created.

We constructed simple 3D photoionization and dust radiative transfer models using the \textsc{mocassin} code (Ercolano et al. 2003, 2005, 2008). The model consists of neutral PAHs and graphitic carbon (optical constants from Li & Draine 2001), and amorphous silicates (optical constants from Laor & Draine 1993) and crystalline forsterite (optical constants from Jager et al. 1998). In both the dust density is proportional to $r^{-2}$ and the grain size distribution is standard Mathis, Rumpl & Nordsieck (1977, hereafter MRN).

We adopted a distance of 1.2 kpc, and a central star luminosity and temperature of $4.25 \times 10^3$ L$_\odot$ and 55 000 K. We varied the shell radii, the position of the silicates and carbon, and the mass of dust in each shell to obtain a good fit to the data. For model to fit the ISO spectrum and the SOFIA data points, the carbon material has to be in the inner parts, while the silicates need to be in the outer parts of the nebula. If silicate dust is present in the inner warmer regions of the shell, emission features at 10 and 18 µm are predicted, opposite to what it is observed. The total mass is constrained to within 10 per cent by the models. In the best fitting model (Fig. 4), the carbon shell extends from 1.5 to 5.5 arcsec on the plane of the sky and contains $6.6 \times 10^{-5}$ M$_\odot$ of dust, while the silicate shell extends from 5.5 to 7.5 arcsec and has a mass of $8.1 \times 10^{-5}$ M$_\odot$.

The carbon dust mass-loss rate, assuming the switch happened 2800 yr ago, would be $2.4 \times 10^{-8}$ M$_\odot$ yr$^{-1}$, so a total mass-loss rate of $7.0 \times 10^{-6}$ M$_\odot$ yr$^{-1}$ if the gas-to-dust ratio is 300. The silicate dust mass-loss rate would be $5.4 \times 10^{-8}$ M$_\odot$ yr$^{-1}$ for a total mass-loss rate of $1.6 \times 10^{-5}$ M$_\odot$ yr$^{-1}$.

The observed C/O ratio in the ionized region is 1.59 (Bernard-Salas et al. 2003). Assuming instant, homogeneous mixing, and a solar O abundance, an initial C/O ratio of 0.8, and an envelope mass of 0.01 M$_\odot$, the numbers suggest that the thermal pulse dredged-up $4 \times 10^{-2}$ M$_\odot$ of carbon, this is almost the same value we obtain. However, the current wind from the central star has C/O = 30 ± 15 (Yu et al. 2009), so that even if the mass-loss rates are low, the current wind can still add significantly to the carbon budget. The range of C/O and the somewhat low carbon dust mass may suggest that the carbon dredge-up was not instantaneous but that the C/O increased while the mass-loss continued.

Independently, a stellar evolutionary model was constructed with an initial stellar mass of 1.5 M$_\odot$ and solar metallicity, fairly typical values for a Galactic disc PN. This model (Fig. 5) derives a peak mass-loss rate at the point where the model star becomes C-rich of $\sim 1 \times 10^{-5}$ M$_\odot$ yr$^{-1}$ and a final central star remnant mass of 0.62 M$_\odot$ (Karlaas 2014), closely matching those values determined from the observations and photoionization modelling. The model shows that the typical interpulse period near the end of the star’s life is $\sim 1 \times 10^3$ years, while the total duration of a thermal pulse, a third dredge-up, and relaxation back to H-shell burning is of the order of 1000 years or more. The numbers fit the duration of the change from oxygen-rich to carbon-rich mass-loss observed with SOFIA, and confirm that this is the likely explanation. At the last thermal pulse, $4.0 \times 10^{-2}$ M$_\odot$ is dredged-up to the envelope. About $1 \times 10^{-3}$ M$_\odot$ of this is $^{12}$C. This is higher than derived from the
observations. Incomplete mixing (Hajduk et al. 2005) can therefore not be ruled out.

The observations and modelling suggest that the thermal pulse took place while the envelope mass was already very low. This is consistent with an AGB final (or fatal) thermal pulse (AFTP) occurring immediately before the star moves off the AGB. An AFTP can lead to both a considerable enrichment with carbon and oxygen and to the dilution of hydrogen, and also explains the current hydrogen-poor nature of the current central star (Blöcker 2001).

The short duration of this phase compared to the long interpulse period suggests that such observations of carbon dredge-up will remain uncommon. BD+30 3639 provides a rare glimpse into a phase of stellar evolution which is crucial to the origin of carbon in the Universe.

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