

# Solar System Tests

## **Do** Rule Out $\frac{1}{R}$ Gravity

Adrienne Erickcek  
Tristan Smith  
Marc Kamionkowski  
*Caltech*

*arXiv: astro-ph/0610483*

# Overview

- $\frac{1}{R}$  Gravity: its structure and motivation
- Scalar-tensor theory: *Ruled Out!*
- Vacuum solution: *NOT Ruled Out!*
- The Solar System solution: *Ruled Out!*



# $\frac{1}{R}$ Gravity: Cosmic Acceleration without Dark Energy

Carroll, Duwuri, Trodden, Turner PRD 70, 043528 (2004)

The Theory:

$$S_{\text{grav}} = \frac{1}{16\pi G} \int d^4x \sqrt{-g} \left( R - \frac{\mu^4}{R} \right)$$

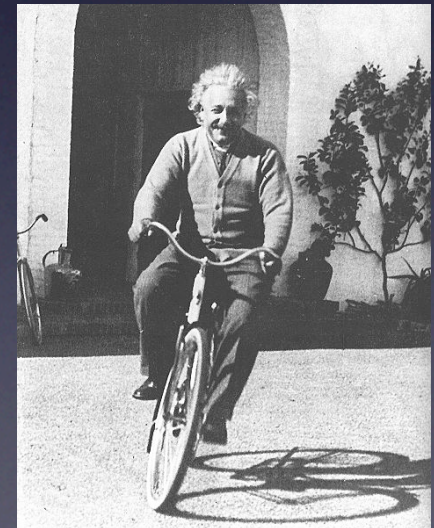
*Action* *Volume* *Mass scale*

The Consequences:

- Constant curvature spacetime is **de Sitter**:

$$\Lambda_{\text{eff}} \sim \mu^2 \sim H_0^2$$

- Late-time acceleration **without dark energy**



# The Scalar-Tensor Twin Theory

*Chiba Phys. Lett. B 575, 1 (2003)*

Consider the **scalar-tensor** theory with the action

$$S_{\text{grav}} = \frac{1}{16\pi G} \int d^4x \sqrt{-g} \left[ \left( 1 + \frac{\mu^4}{\phi^2} \right) \underset{\substack{\uparrow \\ \text{Curvature}}}{R}} - \frac{2\mu^4}{\underset{\substack{\uparrow \\ \text{Scalar Field}}}{\phi}} \right]$$

*Action*      *Volume*

Varying with respect to  $\phi \implies \phi = R$

and then this action is the same as  $\frac{1}{R}$  gravity!



# The Scalar-Tensor Twin Theory

Chiba Phys. Lett. B 575, 1 (2003)

$$S_{\text{grav}} = \frac{1}{16\pi G} \int d^4x \sqrt{-g} \left[ \left( 1 + \frac{\mu^4}{\phi^2} \right) R - \frac{2\mu^4}{\phi} \right]$$

*Action*      *Volume*      *Curvature*      *Scalar Field*

$$S_{\text{grav}} = \frac{1}{16\pi G} \int d^4x \sqrt{-g} [\varphi R - V(\varphi)]$$

**Brans-Dicke theory** with no scalar kinetic term:  $\omega = 0$

Provided that the effective mass of the scalar is small, this theory is **RULED OUT** by Solar System tests.

# Vacuum Solutions

Multamaki and Vilja PRD 74, 064022 (2006)

Field equation for  $\frac{1}{R}$  gravity obtained by varying the metric:

$$8\pi GT_{\mu\nu} = \left(1 + \frac{\mu^4}{R^2}\right) R_{\mu\nu} - \frac{1}{2} \left(1 - \frac{\mu^4}{R^2}\right) Rg_{\mu\nu} + \mu^4 (g_{\mu\nu} \nabla_\alpha \nabla^\alpha - \nabla_\mu \nabla_\nu) R^{-2}$$

*Matter*

*GR-like terms*

*Nasty fourth-order derivatives of the metric!*

Trace Equation: 
$$\frac{8\pi GT}{3} = \nabla_\alpha \nabla^\alpha \frac{\mu^4}{R^2} - \frac{R}{3} + \frac{\mu^4}{R}$$

Assume **vacuum** and **constant curvature**...

**GENERAL RELATIVITY** with a cosmological constant!

$$R = \sqrt{3}\mu^2$$

$$\Lambda_{\text{eff}} = \frac{\sqrt{3}}{4}\mu^2$$



# Vacuum Solutions

*Multamaki and Vilja PRD 74, 064022 (2006)*

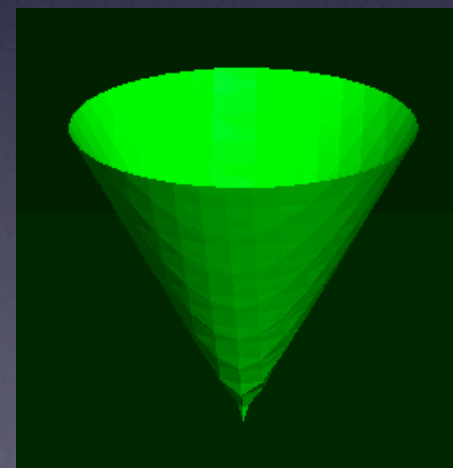
- The **Schwarzschild-de Sitter** metric is a solution to Einstein's equation with a cosmological constant.
- The Ricci scalar  $R$  for this metric is **constant**.
- The **Schwarzschild-de Sitter** metric is a solution to the field equation of  $\frac{1}{R}$  gravity.
- The **Schwarzschild-de Sitter** metric passes all Solar System tests.
- Therefore,  $\frac{1}{R}$  gravity is **NOT RULED OUT!**

# Beware!

This isn't General Relativity...

- In general relativity, the **Schwarzschild-de Sitter** metric is the **only** static spherically-symmetric **solution** to the vacuum equation. (**Birkhoff's Theorem**)
- Birkhoff's Theorem **does not apply** to  $\frac{1}{R}$  gravity.
- Without Birkhoff's Theorem, we have no reason to believe that the **Schwarzschild-de Sitter** metric describes the spacetime in the Solar System.

*Example: A conical spacetime solves Einstein's equations, but it can't be the spacetime around a pressureless string.*





# Finding the Solar System Solution I

In the Sun,  $T \simeq -\rho$

Define a new function:  $c(r) \equiv -\frac{1}{3} + \frac{\mu^4}{R^2(r)}$

Trace of the field eqn:  $\nabla^2 c + \sqrt{3}\mu^2 c = -\frac{8\pi G}{3}\rho$

**The solution:**

$$R = \sqrt{3}\mu^2 \left( 1 - \frac{GM}{r} \right)$$

*Curvature Scalar*

*Background  
value*

*Leading order perturbation  
outside the Sun*

The curvature in the Solar System is not constant!

# Finding the Solar System Solution II

Begin with a **perturbed de Sitter** metric:

$$ds^2 = - [1 + a(r) - H^2 r^2] dt^2 + [1 + b(r) - H^2 r^2]^{-1} dr^2 + r^2 d\Omega^2$$

*Line Element*

Now that we have an expression for  $R$ , the **field equations** are **second-order differential equations** for  $a(r)$  and  $b(r)$ .

*The Linearized Solution*

$$ds^2 = - \left( 1 - \frac{2GM}{r} - H^2 r^2 \right) dt^2 + \left( 1 + \frac{GM}{r} + H^2 r^2 \right) dr^2 + r^2 d\Omega^2$$

*Normal Newtonian  
Term =  $2\Phi$*

*Half of Schwarzschild value:*

$$\gamma = \frac{1}{2}$$



# Conclusions

- Shapiro time delay measurements:

$$\gamma = 1 + (2.1 \pm 2.3) \times 10^{-5}$$

*Bertotti, Iess, Tortora*  
*Nat. 425, 374 (2003)*

- Relating to the equivalent scalar-tensor theory gives the correct PPN parameter for  $\frac{1}{R}$  gravity:  $\gamma = \frac{1}{2}$

- $\frac{1}{R}$  gravity is **RULED OUT** by Solar System tests.

- Analysis may be extended to other modifications of the Einstein-Hilbert action -- stay tuned!